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A refinement of the pseudo crystal structure of scleroclase PbAs₂S₄.[†] By Y. IITAKA and W. No-WACKI, Abteilung für Kristallographie und Strukturlehre, Mineralogisches Institut, Universität Bern, Schweiz

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The crystal structure of scleroclase $PbAs_2S_4$ (from the Lengenbach quarry, Binn Valley, Switzerland) was investigated by Nowacki, Iitaka *et al.* (1961), as a part of the study on sulfosalt minerals (Nowacki & Kunz, 1959, 1961). They noted the occurrence of superlattice reflections corresponding to the $3 \times 1 \times 11$ times greater cell compared with the pseudo-cell. The cell dimensions and the space group of the pseudo-cell were

$$\begin{aligned} a' = & 19 \cdot 62 \pm 0 \cdot 02, \ b' = & 7 \cdot 89 \pm 0 \cdot 01, \ c' = & 4 \cdot 19 \pm 0 \cdot 05 \ \text{\AA}; \\ \beta = & 90^{\circ}, \ C_{2h}^5 = & P2_1/n \ . \end{aligned}$$

They determined the pseudo-structure which is the average structure over the $3 \times 1 \times 11$ such sub-cells, and found that the best agreement between the observed and calculated structure factors for the average structure is obtained when anisotropic temperature factors are introduced for individual atoms. The approximate values of the temperature factor coefficients were determined by difference Fourier syntheses, and the greatest values were observed on the B_{11} and B_{33} of certain atoms, which correspond to the r.m.s. displacement of about 0.3 Å from the mean position along the *a*- and *c*-axis.

It appeared likely that more precise determination of the temperature factor coefficients might indicate the orientation and the magnitude of the modulation which is responsible for the formation of the superstructure. The present communication describes the result of determining more accurate coordinates and temperature factor coefficients by the three-dimensional least-squares method (SFLS programme described by Mills & Rollett 1961, and Cruickshank *et al.*, 1961) which was carried out at the Oxford University Computing Laboratory.

The refinement was started with the atomic coordinates

[†] Structural investigations on sulfosalts from the Lengenbach, Binn Valley (Ct. Wallis). Part 4.—Contribution No. 122.

and anisotropic temperature factor coefficients which were listed in the previous paper (Nowacki, Iitaka *et al.*, 1961). The number of reflections which were used for the present calculation was 797. For the first three cycles the following weighting system was adopted:

$$VW = 1$$
 if $|F_o| < F^*$, $VW = F^*/|F_o|$ if $|F_o| \ge F^*$

 F^* being taken as $2 \cdot 4 F_{o, \min} = 12$. For the 4th cycle this system was modified as

$$VW = |F_o|/F^* \text{ if } |F_o| < F^*, VW = F^*/|F_o| \text{ if } |F_o| \ge F^* \text{ and} VW = F^*/2|F_o| \text{ if } |F_o| \ge 200, (F^* = 12)$$

so as to put smaller weight for the very large $|F_o|$. Altogether four cycles of refinement were carried out;

the process is outlined in Table 1.

Table 1. Th	e refinement	process
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R%	$\Sigma w(F_o - F_c)^2$
18.7	18648
e 17·3	14038
15.2	10892
e 14·7	10085
	$egin{array}{ccc} R\% & & & \ 18\cdot7 & & \ 17\cdot3 & & \ 15\cdot2 & & \ 14\cdot7 & & \ 14\cdot7 & & \ \end{array}$

The average parameter changes during the four cycles of refinement were about 0.03 Å and 0.02 Å in the *a*- and *b*-directions. However, the S_2 atom moved as much as 0.1 Å and 0.04 Å along the *a*- and *b*-axis. At the same time, we noticed a significant increase in the temperature factors which had already been indicated as a residual anisotropy in the difference projections. The final atomic parameters and the standard deviations obtained from the 4th cycle are listed in Table 2.

The interatomic distances are summarized in Table 3.

Table 2. Atomic coordinates and temperature factor coefficients with the standard deviations

	x	y	z	B_{11}	B_{22}	B_{33}	B_{12}
Pb	0.19373 ± 0.00014	0.08717 ± 0.00029	0.25000	$\begin{array}{c} 7 \cdot 09 \\ \pm \ 0 \cdot 13 \end{array}$	$\begin{array}{c} 4.78 \\ \pm 0.07 \end{array}$	$\begin{array}{c} 3 \cdot 39 \\ \pm \ 0 \cdot 07 \end{array}$	-0.40 ± 0.06
As_1	-0.12773 ± 0.00037	-0.50402 ± 0.00079	0.25000	5.12 ± 0.17	4.25 ± 0.16	$\begin{array}{c} 7 \cdot 76 \\ \pm \ 0 \cdot 36 \end{array}$	-0.53 ± 0.14
As_2	0.00459 ± 0.00055	0.80435 ± 0.00130	0.25000	$9 \cdot 24 \\ \pm 0 \cdot 35$	$7.61 \\ \pm 0.30$	$\begin{array}{c} 6\cdot 74 \\ \pm \ 0\cdot 31 \end{array}$	-4.73 ± 0.27
$\mathbf{S_1}$	-0.22125 ± 0.00068	-0.35874 ± 0.00131	0.25000	$\begin{array}{c} 4 \cdot 74 \\ \pm 0 \cdot 33 \end{array}$	$\begin{array}{c} 2 \cdot 82 \\ \pm 0 \cdot 24 \end{array}$	$2 \cdot 92 \\ \pm 0 \cdot 34$	-1.10 ± 0.25
S_2	0.15971 ± 0.00086	0.67467 ± 0.00170	0.25000	$\begin{array}{c} 4 \cdot 11 \\ \pm 0 \cdot 45 \end{array}$	2.70 ± 0.31	$\begin{array}{c} 12 \cdot 75 \\ \pm 1 \cdot 22 \end{array}$	$1 \cdot 27 \\ \pm 0 \cdot 30$
S_3	-0.07790 ± 0.00064	-0.97791 ± 0.00152	0.25000	$2 \cdot 98 \pm 0 \cdot 30$	$\begin{array}{c} 3 \cdot 41 \\ \pm 0 \cdot 30 \end{array}$	$\begin{array}{c} 6 \cdot 12 \\ \pm 0 \cdot 54 \end{array}$	-0.75 ± 0.25
S_4	0.05195 ± 0.00083	0.33385 ± 0.00139	0.25000	3.89 ± 0.43	$2 \cdot 07 \pm 0 \cdot 26$	$\begin{array}{c} 19{\cdot}02 \\ \pm 1{\cdot}62 \end{array}$	-0.28 ± 0.26

Temperature factors take the form, $\exp -\frac{1}{4}(B_{11}h^2a^{*2} + B_{22}k^2b^{*2} + B_{33}l^2c^{*2} + 2B_{12}hka^*b^*)$ assuming the mirror planes exist at $z = \frac{1}{4}$ and $\frac{3}{4}$.

Table 3. Coordination of the Pb and As atoms

Pb:	$2.95 \text{ Å} (S_2) + 3.05(2S_1) + 3.22(2S_1) + 3.22(2S_3)$
	$+3.33(S_2) + 3.42(S_4)$
$As_1:$	$2.14(S_1) + 2.56(2S_2) + 2.92(2S_4)$
As_2 :	$2 \cdot 36(S_3) + 2 \cdot 59(2S_4) + 2 \cdot 87(2S_3)$

It should be noticed that markedly large temperature factors are observed: the B_{11} of Pb, As₂ and the B_{33} of As₁, S₂, S₃, S₄, among them the value of the B_{33} for S₂ and S₄ are surprising. In view of the fact that this structure is the average structure over the 3×11 sub-cells which are extended in the *a*- and *c*-axis direction, it may be considered that these values indicate the direction and a measure of the magnitude of the atomic displacements which would most likely occurr within the superstructure.

Table 4. R.m.s. displacement of the atoms along the three principal crystallographic axes

	a	b	c
\mathbf{Pb}	0·30 Å	0.25 Å	0·21 Å
As_1	0.26	0.23	0.31
As_2	0.34	0.31	0.29
S_1	0.25	0.19	0.19
S_2	0.23	0.19	0.40
S_3	0.19	0.21	0.28
S_4	0.22	0.16	0.49

In Table 4, the root-mean-square displacement of the atoms along the *a*- and *c*-axis are listed which are calculated by the formula $B = 8\pi^2 \overline{\mu^2}$.

It is interesting to note that as illustrated in Fig. 1, the sulfur atoms which have an extraordinary large B_{33} value are those which connect the pyramidal AsS₃-groups to form a chain. If we assume that these sulfur atoms are displaced, say for example, 0.5 Å either upwards or downwards along the *c*-axis, some of the As-S distances would be about 2.17 Å (this is the normal As-S distance) while others would be about 2.98 Å which is too large to assume that a strong bonding exists between the As and S atoms. As discussed in the previous paper (Nowacki, Iitaka *et al.*, 1961), there is a morphological evidence to assume that the As-S-chains do not extend throughout the crystal.

However, as the environment of the Pb atoms is changed by the displacement of the S atoms forming a gap, it might happen that the Pb atoms are displaced (probably along the a-axis) so as to fill up the gap; the complete analysis of the super structure may be very complicated.

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Fig. 1. The chain of the AsS_3 -group found in scleroclase (there are two kinds of chains each constituting $As_1S_1S_2$ and $As_2S_3S_4$).

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